

Isotropic multi-gap superconductivity in $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ from thermal transport and spectroscopic measurements

This content has been downloaded from IOPscience. Please scroll down to see the full text.

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 193.52.23.4

This content was downloaded on 05/01/2015 at 14:11

Please note that [terms and conditions apply](#).

Isotropic multi-gap superconductivity in $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ from thermal transport and spectroscopic measurements

Steven Ziemak¹, K Kirshenbaum¹, S R Saha¹, R Hu¹, J-Ph Reid²,
R Gordon², L Taillefer^{2,3}, D Evtushinsky⁴, S Thirupathaiah⁴, B Büchner⁴,
S V Borisenko⁴, A Ignatov⁵, D Kolchmeyer⁵, G Blumberg⁵ and J Paglione^{1,3}

¹ Center for Nanophysics and Advanced Materials, Department of Physics, University of Maryland, College Park, MD 20742, USA

² Département de Physique & RQMP, Université de Sherbrooke, Sherbrooke, Québec, J1K 2R1, Canada

³ Canadian Institute for Advanced Research, Toronto, M5G 1Z8, Canada

⁴ Institute for Solid State Research, IFW-Dresden, PO Box 270116, D-01171 Dresden, Germany

⁵ Department of Physics and Astronomy, Rutgers, The State University of New Jersey, Piscataway, NJ 08854, USA

E-mail: paglione@umd.edu

Received 24 July 2014, revised 16 September 2014

Accepted for publication 22 September 2014

Published 3 December 2014



CrossMark

Abstract

Thermal conductivity, point contact spectroscopy, angle-resolved photoemission and Raman spectroscopy measurements were performed on $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ single crystals obtained from the same synthesis batch in order to investigate the superconducting energy gap structure using multiple techniques. Low temperature thermal conductivity was measured in the superconducting state as a function of temperature and magnetic field, revealing an absence of quasiparticle excitations in the $T \rightarrow 0$ limit up to 15 T applied magnetic fields. Point-contact Andreev reflection spectroscopy measurements were performed as a function of temperature using the needle-anvil technique, yielding features in the conductance spectra at both 2.5 meV and 7.0 meV scales consistent with a multi-gap scenario. Angle-resolved photoemission spectroscopy probed the electronic band structure above and below the superconducting transition temperature of $T_c = 23$ K, revealing an isotropic gap of magnitude ~ 3 meV on both electron and hole pockets. Finally, Raman spectroscopy was used to probe quasiparticle excitations in multiple channels, showing a threshold energy scale of 3 meV below T_c . Overall, we find strong evidence for an isotropic gap structure with no nodes or deep minima in this system, with a 3 meV magnitude gap consistently observed and a second, larger gap suggested by point-contact spectroscopy measurements. We discuss the implications that the combination of these results reveal about the superconducting order parameter in the $\text{BaFe}_{2-x}\text{Pt}_x\text{As}_2$ doping system and how this relates to similar substituted iron pnictides.

Keywords: iron pnictides, thermal conductivity, point contact spectroscopy, ARPES, Raman spectroscopy

(Some figures may appear in colour only in the online journal)

1. Introduction

Since their discovery [1], the nature of the superconducting gap structure and pairing symmetry of iron-based superconductors (FeSCs) have attracted considerable attention [2].



Content from this work may be used under the terms of the Creative Commons Attribution 3.0 licence. Any further distribution of this work must maintain attribution to the author(s) and the title of the work, journal citation and DOI.

Of particular interest are the 122 iron pnictides, a class of intermetallic compounds with the ThCr_2Si_2 structure in which superconductivity is induced by several methods of chemical substitution. Aside from the more commonly studied substitution series of 3d transition metal elements for iron in $\text{AFe}_{2-x}\text{T}_x\text{As}_2$ ($\text{A} = \text{Ba}, \text{Sr}, \text{Ca}$; $\text{T} = \text{Fe}, \text{Co}, \text{Ni}, \dots$), previous work has shown that 5d transition metal substitution also induces superconductivity, for example in the series $\text{BaFe}_{2-x}\text{Pt}_x\text{As}_2$ where the maximum transition temperature T_c approaches ~ 25 K [3, 4]. Studies to date on this system have focused on crystal structure, resistivity, magnetic susceptibility, and specific heat, but more detailed studies of the pairing energy gap are lacking and would be useful for comparisons to the more heavily studied $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ system.

In understanding the electronic structure of superconductors it is essential to use a comprehensive experimental approach, as each measurement technique has its own unique limitations and opportunities for error. For example, angle-resolved photoemission spectroscopy (ARPES) is considered to be one of the most direct ways to measure the momentum-space structure of the superconducting energy gap Δ , but in many cases ARPES studies on FeSCs have disagreed with bulk probes on the presence or absence of nodes in the order parameter. This apparent paradox may be a consequence of the surface-sensitive nature of the ARPES measurement or experimental limitations, where issues such as surface reconstruction, surface depairing, and/or resolution issues may impair the interpretation of experimental results [5]. Other measurements designed to probe the gap structure similarly suffer from probe-dependent circumstances or are simply less direct probes of a gap's angular dependence. For this reason, we have taken the approach of comparing and contrasting four separate techniques: thermal conductivity, point-contact spectroscopy (PCS), ARPES, and Raman spectroscopy to investigate the $\text{BaFe}_{2-x}\text{Pt}_x\text{As}_2$ system. Furthermore, we utilize samples for each experiment taken from the same synthesis batch, minimizing the error associated with comparisons of experimental results obtained from samples with varying quality, properties and origins, an issue particularly important in systems where chemical substitution is required.

Thermal conductivity measurements provide a powerful probe of electronic excitations in the zero temperature limit, when phonons have been frozen out. In a conducting solid at low temperatures, the temperature dependence of thermal conductivity κ can be modeled as $\kappa(T) = aT + bT^\alpha$, where the T -linear term aT represents the electronic contribution to heat conduction, while the T^α term represents the phonon contribution, where $2 \leq \alpha \leq 3$. By measuring κ/T in the zero temperature limit the electron contribution can be isolated, given by $a \equiv \kappa_0/T$. It has been shown that the presence or absence of this residual thermal conductivity at zero magnetic field and its evolution in field can reveal the presence or absence of nodes or zeroes in the order parameter and provide additional information about gap structure [6]. For example, in the d-wave cuprate superconductors κ_0/T is nonzero at zero magnetic field [7]. The presence of this same residual thermal

conductivity has also been used as evidence for nodal s-wave or d-wave pairing in FeSCs such as $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ [8] and KFe_2As_2 [9]. Likewise, its absence has been taken as evidence for a nodeless s-wave gap structure in the 18 K superconductor LiFeAs [10]. As a function of magnetic field, the evolution of κ_0/T can provide further information about energy scale anisotropies or the presence of multiple gaps with different energy scales, and has been used to differentiate between s-wave superconductors with a single gap (or gaps of equal magnitude on multiple Fermi surface (FS) pockets) and those with multiple gaps of significantly different sizes, e.g., differentiating doubly-gapped NbSe_2 from singly-gapped V_3Si [11].

Another tool for probing electronic excitations is PCS, also known as quasiparticle scattering spectroscopy. Measurements of electrical conductivity versus dc bias voltage across a normal metal/superconductor junction provide an indirect probe of density of states, and fitting such spectra to the Blonder–Tinkham–Klapwijk (BTK) theory [12] allows precise determination of gap size(s) as well as pairing symmetry [13]. PCS measurements on various doped FeSC systems have resulted in widely varying conclusions about their pairing symmetries, at times contradicting the results from other measurement techniques. Point-contact measurements in the Andreev reflection regime have shown evidence for a two-gap structure in $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ [14, 15], while $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ spectra have been fit to both one- and two-gap s-wave models depending on the study and even the fitted features in a single spectrum [13, 14] (i.e., inclusion or exclusion of zero-bias or higher bias features). Finally, PCS measurements of $\text{BaFe}_{2-x}\text{Ni}_x\text{As}_2$ have been used to argue for a two-gap s-wave structure at low Ni concentrations with nodes or deep minima emergent above optimal doping [16].

ARPES provides a more direct picture of pairing symmetry and gap magnitude. In the case of 122 FeSCs, previous work has found two nodeless and nearly isotropic gaps of different sizes on various FS pockets in $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ [17]. ARPES measurements on $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ have shown evidence for hole pockets disappearing upon increasing Co substitution [18], while a separate study focusing on gap structure [19] observed the presence of two isotropic gaps and suggested a connection between FS nesting and Cooper pairing based on comparison with the hole-doped $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ case. Finally, to date less extensive studies of Raman spectroscopy have been completed in the 122 FeSC materials [13, 20], where one study has shown band anisotropy in $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ supporting a nodal s-wave model [21].

In this article, we report the results of four independent studies of the superconducting gap structure in the FeSC material $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ with $T_c = 23$ K. We find consistent observations of an isotropic superconducting gap structure in this material using four techniques. Thermal conductivity measurements show no residual transport in the zero temperature limit under applied magnetic fields of up to 15 T, implying an isotropic, fully gapped s-wave structure with no nodes or deep minima. PCS measurements reveal features in the conductance spectra consistent with a multi-gap scenario.

Fitting spectra to a two-gap BTK theory shows that these features are consistent with the presence of two gaps of magnitude $\Delta_1 = 2.5$ meV and $\Delta_2 = 7.0$ meV. ARPES measurements show a 3 meV energy gap in the density of states that develops below T_c and whose size is independent of direction in k -space, confirming the isotropic nature of the smaller gap structure. Raman spectroscopy reveals excitations in the A_{1g} and B_{2g} channels with thresholds corresponding to a 3 meV gap size consistent with the other measurement techniques.

Details of sample growth and characterization are outlined in section 2, followed by experimental results for thermal conductivity (section 3), PCS (section 4), ARPES (section 5) and Raman spectroscopy (section 6), and general discussion and comparisons in section 7 followed by conclusions in section 8.

2. Sample characterization

Single crystals of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ were grown from prereacted FeAs and PtAs powders and elemental Ba using the FeAs self-flux method [22], which yielded large crystals with typical dimensions $0.1 \times 1 \times 2$ mm³ [3]. Single crystal x-ray diffraction and Rietveld refinement determined a precise Fe : Pt ratio and an exact chemical formula of $\text{BaFe}_{1.906(8)}\text{Pt}_{0.094(8)}\text{As}_2$. Previous x-ray measurements have shown that Pt substitution reduces the c -axis length and c/a ratio while increasing the a -axis length and unit cell volume compared to pure BaFe_2As_2 [3]. Characterization by energy- and wavelength-dispersive x-ray spectroscopy verified the same substitutional concentrations across several specimens.

As shown in figure 1, resistivity measurements exhibit the expected metallic behavior and a sharp transition to the superconducting state with an onset of $T_c = 23$ K and zero resistance achieved by 21.5 K for a transition width of $\Delta T_c < 1.5$ K. No kink is observable at higher temperatures that would indicate a magnetically ordered phase, as is seen in the parent compound and underdoped samples [4, 23]. Magnetic field suppresses T_c down to approximately 16 K in a 14 T field (not shown), consistent with an upper critical field of approximately 45 T as determined by fits to the Werthamer–Helfand–Hohenberg approximation as demonstrated previously [3]. However, $H_{c2}(0)$ may be as high as 65 T based on a linear fit to $H_{c2}(T)$, which would be consistent with similar materials such as $\text{SrFe}_{2-x}\text{Ni}_x\text{As}_2$ [24].

DC magnetic susceptibility was measured with field applied along the ab -plane, under zero-field-cooled conditions and field cooled with a 1 mT applied field. As shown in the inset of figure 1, a sharp transition is observed at $T_c = 23$ K into the diamagnetic state with a full-volume shielding fraction observed, as indicated by $4\pi\chi = -1$ by 17 K. Previous specific heat measurements [3] have confirmed the bulk nature of the transition, with a jump in $C_p(T)$ that occurs slightly below T_c . The size of the jump was estimated as $\Delta C_p/T_c \approx 20$ mJ mol⁻¹K⁻¹, and assuming the

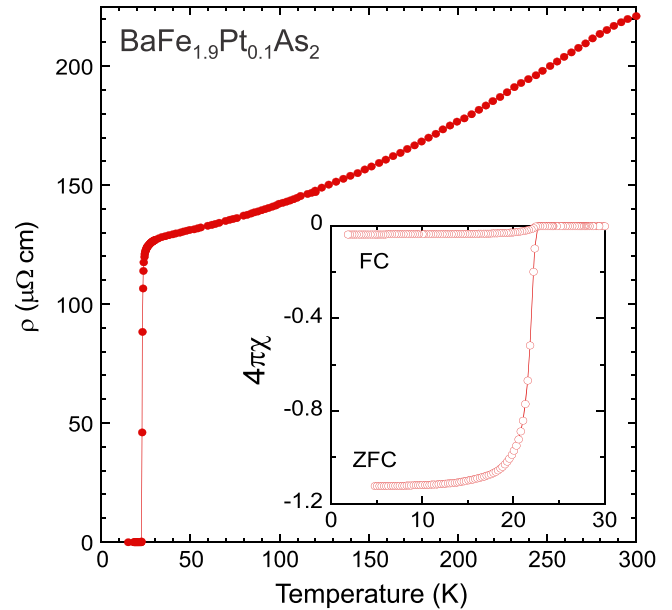


Figure 1. Characteristic properties of the superconducting transition in single-crystal samples of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ with resistivity (main panel) and magnetic susceptibility (inset) measurements exhibiting sharp transition features consistent with a transition temperature at $T_c = 23$ K, as observed in other bulk measurements [3]. Magnetic susceptibility measurements performed in zero-field-cooled (ZFC) and field-cooled (FC) conditions with a magnetic field of 1 mT indicate a full 100% superconducting volume fraction (inset).

BCS weak-coupling ratio of $\Delta C_p/\gamma T_c = 1.43$ applies yields a value for $\gamma = 16(2)$ mJ mol⁻¹K⁻² [3].

3. Thermal conductivity

Thermal conductivity was measured using a one-heater, two-thermometer steady state technique as a function of temperatures down to 60 mK in a dilution refrigerator. Temperature sweeps were repeated in different fixed magnetic field values from 0 T to 15 T, applied in both parallel ($H \parallel ab$) and perpendicular ($H \parallel c$) orientations with respect to the crystallographic basal plane. Measured data, presented as κ/T versus T and shown in figure 2, were fit to the form $\kappa/T = a + bT^{\alpha-1}$ in order to extrapolate to the $T \rightarrow 0$ limit and obtain the residual electronic term κ_0/T .

The residual electronic contribution κ_0/T clearly approaches a value of zero as $T \rightarrow 0$, indicating a lack of zero energy quasiparticles in the superconducting state. (The fitted value of κ_0/T is actually slightly negative, but is within experimental error of zero). As a point of comparison, one can calculate the expected $T \rightarrow 0$ limit electronic contribution to thermal conductivity from the Wiedemann–Franz law, $\kappa_0/T = L_0/\rho_0$, where $L_0 = (\pi^2/3)(k_B/e)^2 = 2.44 \times 10^{-8}$ W Ω K⁻². Using the estimated residual zero-temperature resistivity value of $\rho_0 = 125$ $\mu\Omega$ cm based on extrapolating the zero-field data (see figure 1) yields an approximate value of $L_0/\rho_0 = 0.195$ mW K⁻² cm⁻¹, shown in figure 2 as blue diamonds on

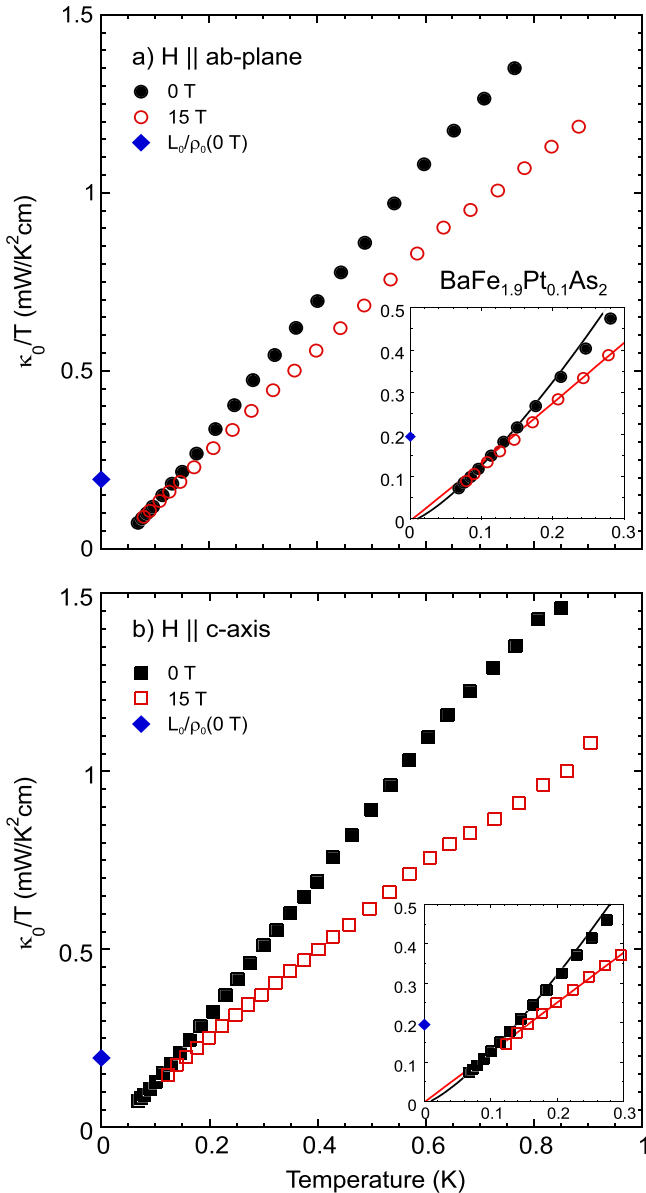


Figure 2. Low temperature thermal conductivity of a crystal of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ measured in zero (black, closed symbols) and 15 T (red, open symbols) applied magnetic fields, oriented along a) the crystallographic ab basal plane (circles) and b) the c -axis direction (squares). Zero field measurements were repeated in each orientation using the same crystal for direct comparison. Insets show a zoom of the low-temperature detail, with solid lines representing power law fits of the form $\kappa/T = a + bT^{\alpha-1}$ to the data points below 150 mK. Blue diamonds in all panels represent the normal state electronic thermal conductivity estimated using the Wiedemann–Franz law (see text).

the $T = 0$ axis. As evident in figure 2, this estimated normal state limit is many orders of magnitude larger than the maximum possible extrapolated value of κ_0/T allowed by error, suggesting a fully gapped superconducting order parameter with no nodes or deep minima. This is to be contrasted with the finite κ_0/T expected for superconductors with either symmetry-imposed or accidental nodes in their gap structure, such as observed in the d-wave superconductor $\text{Tl}_2\text{Ba}_2\text{CuO}_{6+\delta}$ (Tl-

2201) [7] or in c -axis transport measurements of $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ away from optimal doping [8], respectively.

The absence of any low-energy excitations is also evident in magnetic field measurements up to 15 T, or approximately 25% of $H_{c2}(0)$, as shown in figure 2 for both $H \parallel ab$ and $H \parallel c$ orientations in panels (a) and (b), respectively, and summarized as a function of field in figure 3. As a function of increasing magnetic field, it has been shown that single-gap isotropic s-wave superconductors such as Nb [25] continue to lack low-energy quasiparticle excitations with increasing magnetic field until vortex core bound states begin to delocalize, a process that proceeds exponentially with the ratio of intervortex spacing to coherence length (or essentially the magnitude of magnetic field) up to the upper critical field. Multiband s-wave superconductors such as NbSe_2 also lack any κ_0/T term in zero field due to lack of nodes in the gap structure. However, the vortex core delocalization process can be accelerated in such cases due to the presence of reduced gap magnitudes on one or more bands, resulting in an onset of low-energy excitations and a finite and increasing value of κ_0/T at fields much smaller than H_{c2} [11].

The evolution of κ_0/T with magnetic field in $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ is compared to the behavior of several characteristic superconductors in figure 3(a), and to varying behaviors observed in the FeSC family in figure 3(b). The lack of any increase in residual term in $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ is made even more clear when plotted as a function of field normalized to $H_{c2}(0)$ (using 65 T for $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ as a conservative estimate): the evolution in $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ is comparable to that of elemental Nb [25], and is in stark contrast to that of the d-wave and multiband s-wave superconductors Tl-2201 and NbSe_2 , respectively. Within the FeSC family, such a flat response is not unprecedented, and is in fact favorably comparable to the full isotropic gap scenarios deduced for both $\text{Ba}_{0.6}\text{K}_{0.4}\text{Fe}_2\text{As}_2$ [26] and LiFeAs [10].

In the case of a multi-band scenario, which is the case for $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ as measured by ARPES (see section 5), the suppression of the maximum energy gap on portions of the band structure can lead to an enhancement of tunneling between vortex core bound states, resulting in an effectively reduced critical field for those portions. In the case of NbSe_2 , this field scale is identified with a shoulder in $\kappa_0/\kappa_N(H)$ near $H^* \sim H_{c2}/9$, consistent with an energy gap ratio $\Delta_{\min}/\Delta_{\max} \simeq 1/3$ as follows from the fact that $H_{c2} \propto \Delta^2/v_F^2$, where v_F is the Fermi velocity [11]. In comparison, in $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ the lack of any rise whatsoever in $\kappa_0/\kappa_N(H)$ in the field range studied up to $H_{c2}/4$ suggests a much weaker reduction from the maximum gap magnitude on those portions of the FS with a smaller gap, assuming the same assumptions apply. However, as discussed for the case of LiFeAs [10], the lack of any signature of a rise in $\kappa_0/\kappa_N(H)$ cannot completely rule out the existence of a smaller gap in the band structure, whose presence may be compatible these observations if for instance the conductivity contribution of the small-gap FS is much smaller than that of the large-gap component, or if coherence lengths are comparable due to

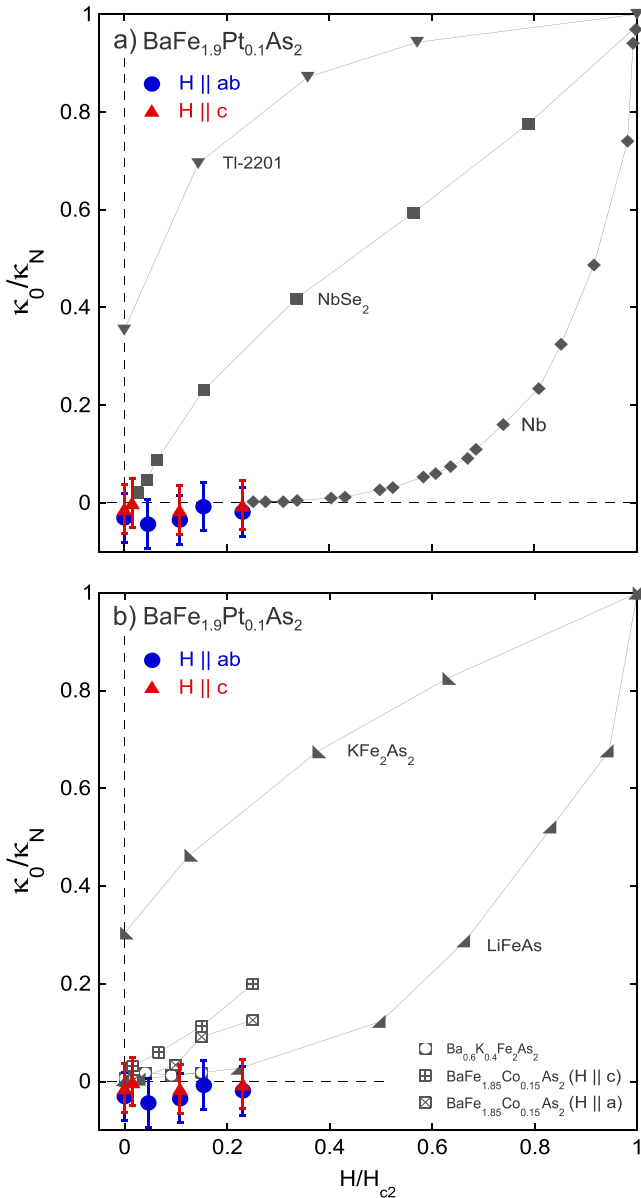


Figure 3. Residual thermal conductivity κ_0/T of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ extracted from temperature sweep measurements and normalized to the normal state conductivity κ_N/T , plotted as a function of reduced magnetic field and compared to several characteristic superconductors. As shown in panel (a), the lack of increase in κ_0/κ_N in $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ is comparable to the behavior of isotropic, single-gapped *s*-wave superconductors such as elemental Nb [25] that exhibit an exponentially slow activated increase with field, and is in contrast to the cases for multi-band *s*-wave superconductor NbSe_2 [11] and *d*-wave nodal superconductor $\text{Tl}_2\text{Ba}_2\text{CuO}_{6+\delta}$ [7], which exhibit a much faster rise with field due to nodal or small-gap low energy quasiparticle excitations. Panel (b) compares $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ to the range of behaviors found in iron-based superconductors such as KFe_2As_2 [26], $\text{Ba}_{0.6}\text{K}_{0.4}\text{Fe}_2\text{As}_2$ [26], $\text{BaFe}_{1.85}\text{Co}_{0.15}\text{As}_2$ [8], and LiFeAs [10] as noted.

scaling of Fermi velocities (i.e., since $\xi_0 \sim v_F/\Delta$). In fact, the case of LiFeAs , as will be discussed below, is a pertinent comparison for the multi-band (and multi-gap) nature of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ because recent ARPES experiments on LiFeAs do indeed find indications of two gap magnitudes of

approximately 2–3 meV and 5–6 meV [27–29], with mild anisotropies reported but with gap minima not reaching below ~ 2 meV. Therefore, it is likely that $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ also shares the situation of having similar coherence lengths across bands, making determination of the gap anisotropies with thermal conductivity difficult and complementary experiments necessary for further understanding.

4. PCS

Point-contact Andreev reflection spectroscopy (PCS) was used to further investigate the gap structure. Needle-anvil junctions were prepared using a sharpened Au tip attached to a specialized probe with a moveable stage, which was used to create contact between the tip and the sample. The samples used were large platelets extracted from FeAs flux and were freshly cleaved before applying contacts and inserting them into the probe. The junctions whose results are presented below had typical dc junction resistances on the order of 10–20 Ω . ac resistance was measured against dc bias voltage with a lock-in amplifier used to stabilize the ac signal, over a range of temperatures using a ^4He dipper cryostat.

In a normal metal-superconductor junction, measurements of the current–voltage (*I*–*V*) relation can reveal information about the nature of the superconducting gap(s) if the size of the junction is small compared to the superconductor’s mean free path. Electrons incident on the junction from the normal metal with an energy of less than Δ can either be reflected back or generate a Cooper pair in the superconductor and reflect a hole back into the normal metal to conserve charge, spin, and momentum in a process known as Andreev reflection [30]. In a measurement of conductivity dI/dV versus bias voltage across the junction, the Andreev reflection process yields an enhancement of conductance when the chemical potential difference is less than Δ . For the case of complete Andreev reflection in an isotropic *s*-wave superconductor, this enhancement occurs as a nominal doubling of the conductance that abruptly onsets at the gap energy. In the case of a junction with a large potential barrier between the normal metal and superconductor, conductivity is instead suppressed due to enhanced reflection at the junction, yielding the typical tunneling conductance spectrum.

For an isotropic *s*-wave superconductor, Andreev reflection leads to a flat doubling of the conductance from zero bias up to $\pm\Delta$ [12, 30], while in other cases the conductance can exhibit sharper and larger enhancements near or at zero bias, in particular in nodal superconductors which exhibit sharper, non-flat peaks in their dI/dV spectra. However, it is known that zero-bias conductance peaks can be observed in point contact measurements due to a wide variety of phenomena. Thus, consistency and replicability are crucial in verifying that observed dI/dV features are indeed intrinsic and not artifacts of a measurement problem.

We have measured conductance curves using junctions created with several different $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ crystals obtained from the same synthesis batch. A sampling of these dI/dV curves are shown in figure 4, presented as raw data without any

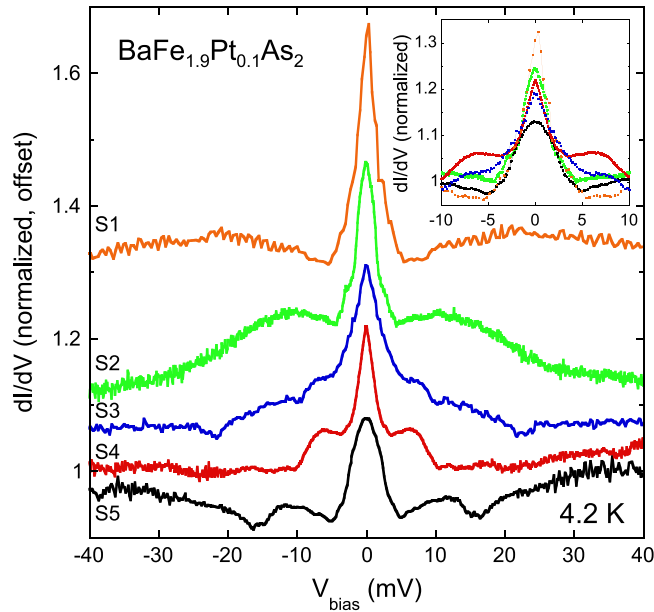


Figure 4. Conductance curves obtained from PCS measurements performed on five $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ samples from the same synthesis batch, offset vertically for clarity. Each has a strong zero bias enhancement of similar width and either a depression or secondary enhancement at higher bias. These curves represent raw data without background subtraction, and each curve is normalized to their high bias values. The inset presents the conductance curves for all five samples plotted without an offset in order to show the overlap of the gross gap features. The green curve (S2) corresponds to the data analyzed in figures 5 and 6.

background normalization (see below) but only normalized to their individual high-bias values (i.e., at 40 meV for each curve) in order to account for differences in junction resistance. Disregarding the slight differences in the shape of the curves, which results from differing background or scattering (γ) contributions, each of the dI/dV curves in figure 4 features an abrupt low-bias peak with a width of approximately 2–3 meV, consistent with Andreev reflection in the superconducting state. Furthermore, all curves exhibit depressions or enhancements in conductance at higher energies (i.e., closer to 5–6 meV). This overlap is emphasized by the inset of figure 4, which presents the same data without vertical offsets.

Figure 5 presents the temperature evolution of PCS conductance spectra for sample S2, now normalizing the dI/dV data to the normal state spectra to remove the background contribution to conductance (see figure 4). This is done by dividing out a polynomial fit to the data measured at 18 K from each temperature data set. Upon cooling below T_c ,⁶ the features described above clearly emerge and evolve with decreasing temperature to reveal a sharp conductance enhancement and a depression at higher bias, suggesting

⁶ Features in the conductance spectra associated with the onset of superconductivity in sample S2 appear to emerge only below 18 K, which is lower than the nominal T_c value of 23 K measured by resistivity, susceptibility and heat capacity [3]. This may be due to critical current effects or possibly a departure from the diffusive Sharvin regime at higher temperatures [31].

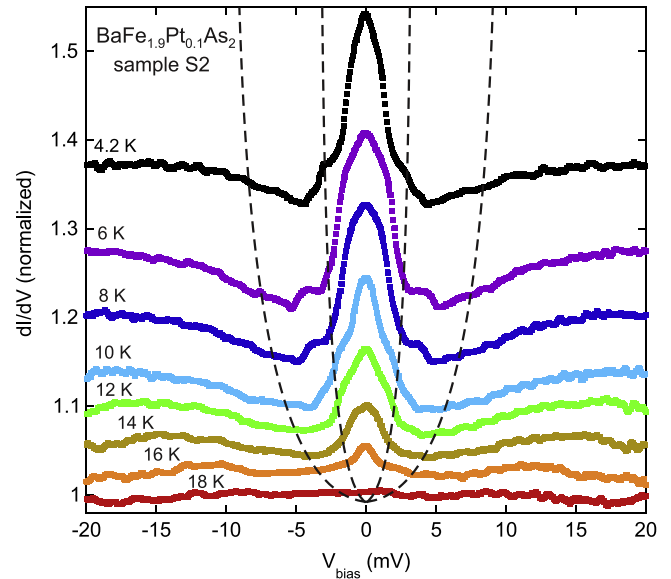


Figure 5. Normalized dI/dV spectra for single crystal $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ with high temperature background conductance divided out over a range of temperatures. Curves are offset for clarity. Dashed lines serve as a guide to the eye to demonstrate that the two superconducting gap features, the low-bias conductance enhancement and higher bias wells, decrease as expected up to T_c (see text for details).

features with at least two energy scales evolving in an order parameter-like fashion.

To better understand the features observed in the conductance spectra, we employ fitting to the BTK model, which includes several parameters which can be used to fit the experimental data [12]. The I - V characteristics of a normal metal-superconductor (N-S) junction are determined by the gap size Δ , a unitless barrier strength Z , and an inelastic scattering energy γ . The value of Z determines the degree of tunneling across the junction: $Z = 0$ indicates complete Andreev reflection, while $Z \rightarrow \infty$ indicates the tunneling limit. In the low- Z case, the dI/dV versus V_{bias} curve contains a peak at zero bias extending out to bias voltages corresponding to approximately $+\Delta$ and $-\Delta$, whereas for large Z , dI/dV is suppressed at zero bias with enhancements at $\pm\Delta$ due to a tunneling gap. The γ term generally changes the shape of the dI/dV curves and is therefore also critical in obtaining a good fit to measured data. In the case of a two-gap BTK model, the fit parameters include two gap sizes, Δ_1 and Δ_2 , two barrier strengths, Z_1 and Z_2 , two scattering energies, γ_1 and γ_2 , and an additional weight factor w , which indicates the proportion of electrons incident on the barrier that interact with either gap. The composite two-gap BTK curve is simply a weighted average of two single-gap curves.

Fitting the conductance spectra described above yields a good fit to the isotropic two-gap BTK formalism, as shown in figure 6 for sample S2 at 4.2 K. In particular, the fit captures both low- and high-bias features of enhancement and depression of conductance, respectively, via two sets of parameters: the low-bias conductance enhancement of this sample can be modeled by the existence of a nodeless gap

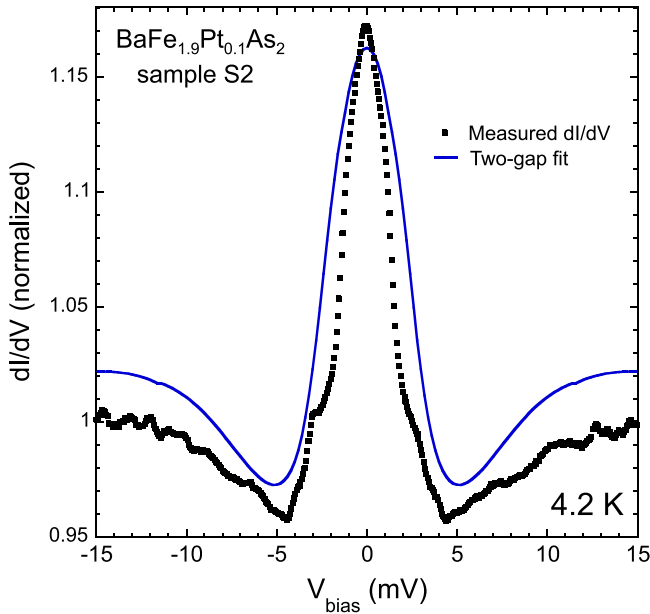


Figure 6. Point contact conductance spectrum of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ sample S2, measured at 4.2 K and normalized to remove the normal state background conductance (see text). Data are fit to an isotropic s-wave two-gap Blonder–Tinkham–Klapwijk model (see text for details), yielding two sets of parameters for a small-gap spectrum ($\Delta_1 = 2.5$ meV, $Z_1 = 0$, $\gamma_1 = 1.0$ meV) and a large-gap spectrum ($\Delta_2 = 7.0$ meV, $Z_2 = 7.0$, $\gamma_2 = 7.0$ meV), with a weighting factor of $w = 0.6$.

with energy scale $\Delta_1 = 2.5$ meV, while the high bias depression is well modelled by including a second gap of magnitude $\Delta_2 = 7.0$ meV such that the fit closely matches the observed spectrum. These features are strikingly similar to those observed in PCS measurements performed on the two-gap superconductor MgB_2 , which also exhibit peak-dip features that were understood in terms of a two-gap scenario with similar gap energies (but of course a much higher T_c of 40 K) [32].

Interestingly, significant differences in the fit parameters in terms of the Z parameter and broadening term γ were necessary to fit the PCS spectra for $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$. For instance, the parameter set for the smaller gap spectrum ($\Delta_1 = 2.5$ meV, $Z_1 = 0$, $\gamma_1 = 1.0$ meV) indicates complete Andreev reflection and a small broadening term, which is typical for such a transparent gap. However the larger gap spectrum ($\Delta_2 = 7.0$ meV, $Z_2 = 7.0$, $\gamma_2 = 7.0$ meV) appears closer to the tunneling limit with a larger Z and exhibits much greater broadening. The weight factor of $w = 0.6$ indicates comparable amounts of Andreev reflection and tunneling into each gap with more conduction into the smaller gap.

While it is not clear why such significant differences were observed in the modeled barrier heights and scattering strengths of the two gaps, such differences are not unprecedented. For example, a two-gap model whose larger gap also has a larger Z was observed in $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ [14]. The unusually large scattering strength deduced by the BTK fit for the larger gap, with $\gamma_2 \simeq \Delta_2$, is also unusual. However, it is in fact not unreasonable considering the high quasiparticle

scattering rates deduced from pair-breaking experiments [10, 33]. Kirshenbaum *et al* investigated pair-breaking scattering in a range of 122 iron pnictide materials with transition metal substitutions at the Fe site [33]. Using their observation of universal pair-breaking behavior in these compounds, one can estimate the transport scattering rate for $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ based on a T_c value of 23 K, yielding $\Gamma \simeq 2.5 \times 10^{13} \text{ s}^{-1}$. This translates to an energy scale of $\hbar\Gamma \approx 16$ meV, which is actually higher but of the same order as γ_2 , suggesting there is a more profound question about strong scattering in the FeSC materials as addressed elsewhere [33]. Alternatively, the contrast between $Z_1 = 0$ and $Z_2 = 7.0$ can find an explanation in terms of a relatively larger ‘impedance mismatch’ between the tip and the large-gap band than for the small-gap band. Following the argument from thermal conductivity measurements on LiFeAs [10] in regard to a mismatch in Fermi velocities, it is possible that the observation of this large contrast in Z values and the (lack of) evidence for multi-band effects in thermal conductivity are not inconsistent.

Finally, the appearance of spectral features at energy scales much larger than the predicted weak-coupling BCS gap value for a 23 K superconductor, $\Delta_{\text{BCS}} = 1.76k_B T_c = 3.5$ meV, is puzzling but also not unprecedented in tunneling and PCS studies. For normal-superconductor junctions outside the ballistic limit, it has been shown that sharp dips can occur at higher bias voltage [34]. A conservative estimate gives an electron mean free path on the order of nanometers [35], which would place our micron-size junctions (assuming a perfect and full contact) in the diffusive rather than ballistic regime, and outside the Sharvin limit. However, as evidenced in figure 4 all sample spectra exhibit features at higher bias voltages, including both gap-like dips and Andreev-like shoulders. It is true that many PCS experiments observe dominant spectral features in a lower-energy range consistent with our prominent features near Δ_1 , such as in $\text{BaFe}_{1.86}\text{Co}_{0.14}\text{As}_2$ [14] and in several 122’s that exhibit a universal $2\Delta/k_B T_c$ ratio of 3.1 [36]. In the case of LiFeAs , observations of a small gap of 1.6 meV ($2\Delta/k_B T_c = 2.2$) have been reported from PCS measurements using Pb and Au tips [37], which do not observe features associated with any larger gap as seen in ARPES experiments [28, 29]. However, many studies have also revealed prominent two-gap features, such as shown in $\text{Ba}_{0.55}\text{K}_{0.45}\text{Fe}_2\text{As}_2$ [14, 15] and $\text{BaFe}_{2-x}\text{Ni}_x\text{As}_2$ up to $x = 0.10$ [16]. Furthermore, a wide range of scanning tunneling spectroscopy experiments have observed superconducting gap magnitudes indicative of strong coupling, with $2\Delta/k_B T_c$ ratios far above the BCS weak-coupling expectation of 3.5 [37]. For example, experiments on optimally doped $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ have observed a superconducting tunneling gap with coherence peaks corresponding to a much larger (average) single-gap value of $\Delta = 6.25$ meV, corresponding to $2\Delta/k_B T_c = 5.73$ [39].

5. ARPES

ARPES allows for direct visualization of electronic structure of materials and precise characterization of the gapping of occupied states by the superconducting order parameter, and

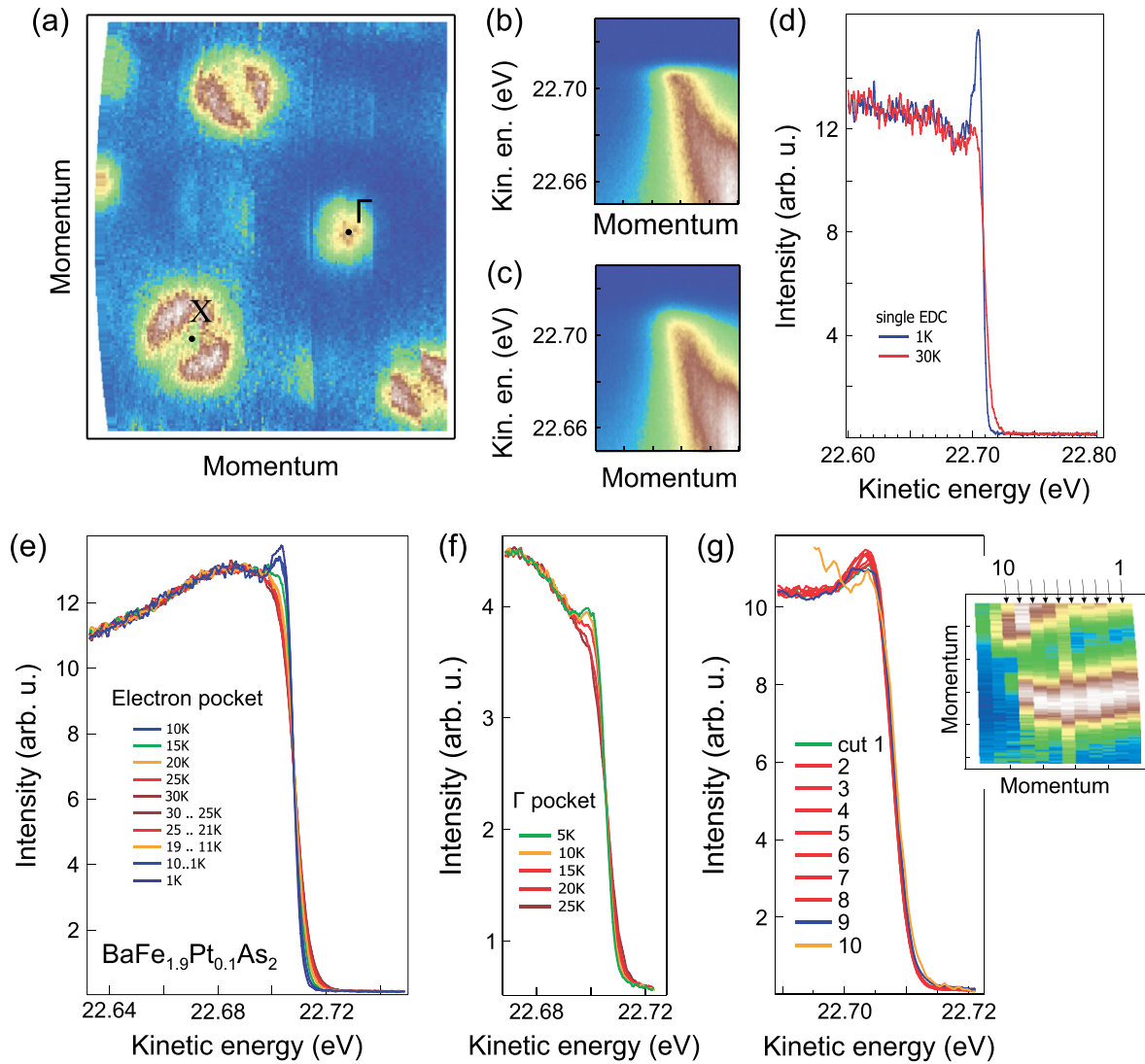


Figure 7. ARPES measurements of $\text{BaFe}_{2-x}\text{Pt}_x\text{As}_2$. (a) Fermi surface map, revealing large electron-like pockets around the Brillouin zone corner and small hole-like pockets at the center. (b) Energy–momentum cut, passing through the electron pocket recorded below T_c . (c) Same cut, recorded above T_c . (d) Single energy distribution curve (EDC) recorded above and below T_c shows appearance of a sharp peak in the superconducting state. (e) Temperature dependence of the integrated EDC (IEDC) recorded from electron pocket for heating up to 30 K and cooling down to 1 K. (f) Temperature dependence of the IEDC for hole pocket. (g) IEDC recorded from different parts of the electron pocket. Positions of cuts are indicated in the mini-map in inset.

is therefore well-suited for identifying and characterizing the momentum-resolved energy scales of the gap function $\Delta(k)$, in particular for elucidating the gap structure of multi-band superconductors.

Measurements were performed at 1^3 end station at BESSY II (Helmholtz-Zentrum für Materialien und Energie). Data were recorded from freshly cleaved samples. Samples were cleaved in high vacuum at low temperatures, exposing mirror-like smooth surfaces. The FSmap of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$, shown in figure 7(a), exhibits large electron pockets at the X point and smaller hole pockets at the Γ point, which is consistent with other electron-doped 122 FeSCs [40]. Energy–momentum cuts presented in figures 7(b) and (c) show occupied bands above and below T_c . Several energy distribution curves (EDCs) were measured and are presented in

figures 7(d)–(g). These curves show the consistent appearance of a peak, which is consistent with a gap opening below approximately 17 K on both the electron and hole pockets. For example, upon cooling below T_c the appearance of a sharp coherence peak and gapping at the Fermi energy are clearly seen in the EDC curve presented in figure 7(d). The width of these features corresponds to approximately a 3 meV gap, which is consistent with the small gap size Δ_1 extracted from fits of the PCS data in section 4.

Integrated energy distribution curves (IEDCs) were also recorded across several cuts at different positions along the electron pocket, as shown in figure 7(g). Each IEDC has a peak of approximately the same width, confirming that the observed electron pocket gap is isotropic. This is comparable to several other 122 FeSCs studied previously, such as

optimally doped $\text{Ba}_{1-x}\text{K}_x\text{Fe}_2\text{As}_2$ [41], $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ [19] and LiFeAs [27], which all show multiple gap sizes with little or no anisotropies.

6. Raman spectroscopy

The polarization-resolved Raman scattering traces electronic density-density correlations driven by the incident and scattered photons of the chosen polarization. In the effective mass approximation for multiband system, Raman coupling to electronic excitations is proportional to $[(m/\hbar^2)\sum_{\alpha,\beta}e_{\alpha}^1e_{\beta}^S\sum_{b,\vec{k}}(\partial^2\epsilon_{\vec{k}}^b/\partial k_{\alpha}\partial k_{\beta})]^2$, with m denoting the electron mass and $\hat{n}_{b\vec{k}}$ the occupation of the Bloch state at momentum \vec{k} in the band b with the dispersion $\epsilon_{\vec{k}}^b$. For the 122 pnictides the two photon fields with cross polarization couple to B_{2g} susceptibility if polarization of light is aligned along a and b crystallographic directions. This scattering geometry mainly samples excitations in the vicinity of the electron pockets. The coupling to the B_{1g} susceptibility has been shown to be weak [42, 43].

In this work Raman spectra were excited with 647 and 476 nm Kr^+ laser lines. Incident laser power between 1.2 and 12 mW was focused into a spot of $50 \times 100 \mu\text{m}^2$ on the freshly cleaved ab -plane crystal surface. The scattered light collected close to the backscattered geometry was focused onto $100 \times 240 \mu\text{m}^2$ entrance slit of a custom triple-stage spectrometer equipped with 1800mm^{-1} holographic gratings. We employed both linearly- and circularly-polarized light to record symmetry-resolved Raman spectra of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$. Data were collected from fourteen samples at 3 and 30 K. The estimated local heating in the laser spot did not exceed 4 K for laser power less than 2 mW.

Secondary emission which depicts B_{2g} symmetry channel is shown in figure 8. In the normal state the response is essentially flat (red line). The response in the superconducting state exhibits a broad peak about $80\text{--}90 \text{cm}^{-1}$ and a threshold at $\approx 45 \text{cm}^{-1}$ (blue line). It flattens out below $\approx 25 \text{cm}^{-1}$. The flat background (uppermost black dashed line) is likely due to second phases in the sample. All 14 samples measured in this Raman study showed nonvanishing backgrounds which amplitude depend on spot position within the cleave.

Figure 8 illustrates the signal decomposition from the mixed phase of doped iron pnictides and highlights the contribution of the major superconducting phase to the Raman response. The embedded flux gives rise to the featureless luminescence signal. The magnitude of this background luminescence (bottom arrow) is evaluated as the minimal contribution for each measured polarization. The Raman response is calculated as $\chi''(\omega) = [d\sigma(\omega)/d\omega - BG]/[n(\omega, T) + 1]$, where $n(\omega)$ is the Bose factor. The resulting response is further decomposed into a sum of contributions from the major superconducting phase and impurity phase(s). These potential impurity phases may not be superconducting or may have a much lower T_c . Based on the flattened section of the experimental data (from the cutoff near 10cm^{-1} to roughly 25cm^{-1}) at 3 K and general flattened response at 30 K, we can assume

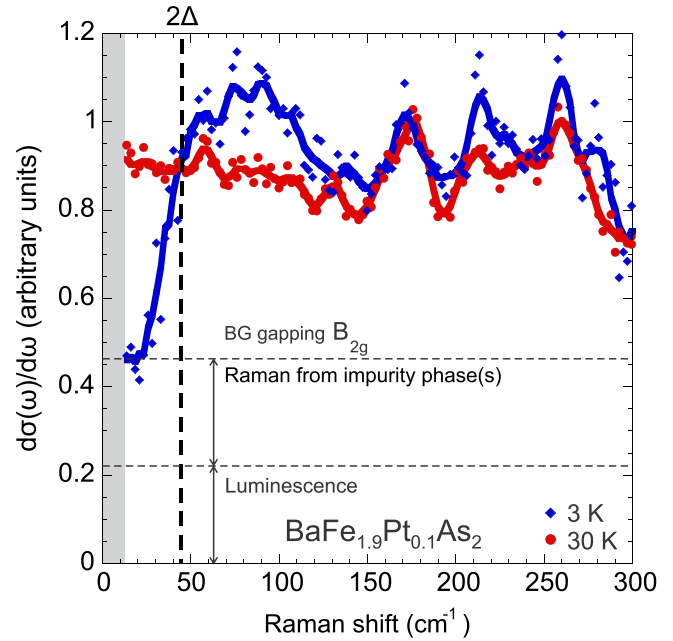


Figure 8. The secondary emission measured from sample S3 with $\lambda_L = 467 \text{nm}$ laser excitation containing Raman and luminescence contributions. Blue diamonds represent data at 3 K, while red circles represent 30 K data. The superconducting gap is indicated by the vertical dashed line labeled 2Δ . Note that the gray area indicates energies below the cutoff value of the spectrometer.

that contributions from any impurity phase(s) can be well approximated by a constant.

Raman response in B_{2g} and A_{1g} symmetry channels is depicted in figures 9(a) and (b), respectively. In the superconducting state data in B_{2g} channel exhibit a threshold at about 45cm^{-1} which is consistent with the fundamental gap magnitude observed by ARPES and PCS. The value of $\Delta = 3 \pm 0.3 \text{meV}$ was confirmed for multiple $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ samples.

7. Discussion

In order to understand these results in context it is instructive to recall the conclusions drawn from similar experiments on Co-substituted BaFe_2As_2 . Results from multiple studies on optimally doped $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ tend to confirm a similar nodeless gap structure. Comprehensive analysis of the phase diagram by Reid *et al* [8] looked at the evolution of gap structure with increasing Co concentration via thermal conductivity. Based on the presence or absence of residual thermal conductivity κ_0/T and its evolution in field, it was shown that line nodes are present in under- and overdoped $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ but that optimally doped $x = 0.148$ samples were nodeless.

Daghero *et al* measured several PCS spectra for overdoped samples with $x = 0.2$ [13]. Current was injected in the ab -plane and parallel to the c -direction using the soft point contact method, in which a small amount of conductive paint is used to create the N-S junction rather than pressing a

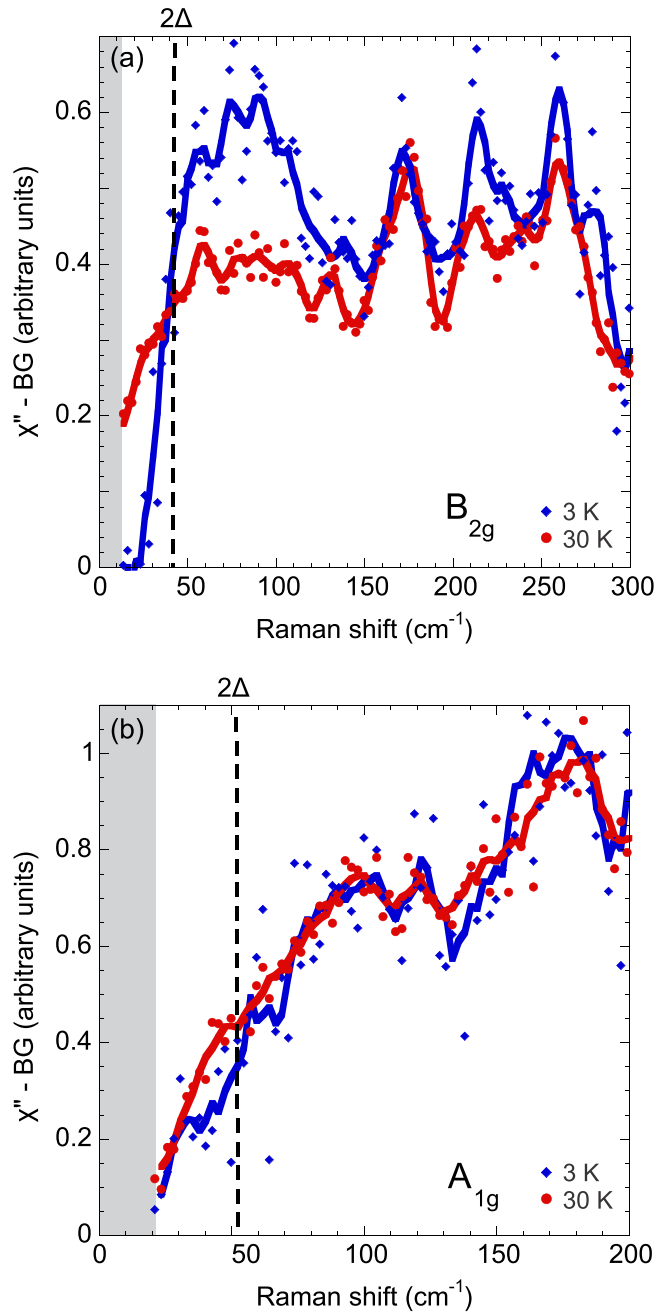


Figure 9. Electronic Raman response (background removed) of sample S3, obtained with 476 nm laser excitation in (a) B_{2g} and (b) A_{1g} channels. Blue (red) dots show superconducting (normal state) response averaged over four (two) experimental scans. Both thick lines represent five-point smoothed data. Black dashed lines mark the energies corresponding to the superconducting gap observed, 2Δ . The gray area again indicates areas below the cutoff energy of the spectrometer.

sharpened tip into the SC sample. Attempts to fit to BTK theory showed a better fit for a two-gap isotropic s-wave model than for a single nodeless gap. These fit curves included estimated gap values of $\Delta_1 = 3.8\text{--}5.2\text{ meV}$ and $\Delta_2 = 8.2\text{--}10.9\text{ meV}$. The authors also reported features in the measured dI/dV curves at higher bias voltage that deviated from the BTK fit which were attributed to strong electron-

boson coupling. PCS measurements were also performed by Samuely *et al* with sharpened Pt tips pressed into freshly cleaved optimally doped ($x = 0.14$) crystals [14]. BTK fits to the measured dI/dV curves showed only a single isotropic gap of magnitude $\Delta = 5\text{--}6\text{ meV}$ suggesting that if multiple gaps are present they are close together in magnitude. Given the significant differences in gap structure observed through thermal conductivity across the $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ phase diagram, it is possible that the stark difference in these results may be attributed to their different Co concentrations.

Terashima *et al* examined the gap structure of $x = 0.15$ samples using ARPES and found strong evidence for two isotropic gaps [19]. EDCs showed two gaps with $\Delta_1 = 5 \pm 1\text{ meV}$ on the electron pocket and $\Delta_2 = 7 \pm 1\text{ meV}$ on the hole pocket. EDCs measured over a range of angles shows that both gaps are isotropic within error.

Muschler *et al* performed Raman spectroscopy measurements on $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ samples with optimal doping ($x = 0.122$) and slight overdoping ($x = 0.17$) [21]. Based on the strong low-temperature shift in the B_{2g} spectrum which varies as $\sqrt{\Delta}$ the authors propose an s-wave state with accidental nodes in the optimally doped sample. Slight overdoping resulted in a B_{2g} peak with greatly diminished amplitude suggesting that the gap is strongly effected by doping and disorder (and potentially sample quality).

The most agreeable conclusion regarding the gap structure of $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ seems to be that of an s-wave gap which is nodeless at optimal doping with nodes appearing at higher and lower Co concentrations which do not appear to be imposed by symmetry and which are present away from the ab -plane. However, some results appear to contradict each other. For example, PCS studies consistently argue in favor of an isotropic gap while Raman spectra at multiple dopings suggest the presence of nodes. Some have suggested reconstruction of the gap at the surface as a means to explain differences between measurements, but this alone does not reconcile disagreements between thermal conductivity, PCS, and Raman spectroscopy, all of which are bulk probes. Another possible explanation is differences in crystal quality, and thus impurity scattering, which has been proposed by Muschler and others.

Compared to the Co-substituted system, our results on $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ offer a more comprehensive conclusion of an isotropic gap structure as all of the samples used were made in the same batch and therefore are expected to have similar impurity concentrations and defect structures. Furthermore, our results all support the conclusion of at least one gap with a consistent magnitude of approximately 3 meV, and all experiments point to an isotropic s-wave gap structure.

Looking at the Co- and Pt-doping comparison more directly, both optimally doped materials lack a residual thermal conductivity at zero field, indicating a fully-gapped superconducting order parameter. For the Co compound, a nonzero κ_0/T emerges at small magnetic fields, while for the Pt compound it remains zero at all fields observed. While the thermal conductivity of each material could only be observed up to a small fraction of H_{c2} , this may suggest subtle

differences in gap morphology. PCS measurements on both yielded BTK fits to nodeless s-wave models. Results from our study with optimal Pt doping seem to compare more directly with those of overdoped $\text{BaFe}_{2-x}\text{Co}_x\text{As}_2$ rather than optimally doped. Both point to a two-gap structure with features present in the spectrum that do not perfectly match the BTK fit, which may indicate strong electron-boson coupling in both materials. ARPES measurements on optimally doped samples in both systems showed no variation in gap magnitude as a function of angle [19]. In the Co-doped compound, two gaps were consistently seen, while only one was observed in our study of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$. Finally, Raman spectra on both compounds feature low temperature enhancements in the A_{1g} and B_{2g} channels at positions which agreed with other reported values of gap size.

While we found no evidence of nodes for optimally doped $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$, under- and overdoped compounds in the $\text{BaFe}_{2-x}\text{Pt}_x\text{As}_2$ system have yet to be explored as far as gap structure is concerned. It has yet to be seen whether similar nodes would appear or if the Pt-substituted system would exhibit significant differences from Co substitution. Therefore, future investigation of the entire $\text{BaFe}_{2-x}\text{Pt}_x\text{As}_2$ phase diagram is certainly warranted.

Our results on $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ can also be used to draw interesting parallels with the multiband superconductor LiFeAs . Both compounds present no evidence of quasi-particle excitations in thermal conductivity measurements at low magnetic fields, and show minimal gap anisotropy in gap structure as determined by ARPES measurements. Both also exhibit comparable small and large gap magnitudes as extracted from PCS measurements and ARPES measurements in the case of LiFeAs , with a $2\Delta/k_B T_c$ ratio of approximately 2 for the smaller gap in each system. The multi-gap nature of both materials would seem to be in disagreement with the thermal conductivity observations, but can be reconciled if one assumes the coherence length is similar between the two gaps. In our PCS observations of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$, the large discontinuity between Z values for the two gaps is consistent with large differences between Fermi velocities, which means that v_F/Δ could be comparable for the two gaps as suggested for LiFeAs [10].

8. Conclusion

In conclusion, we have used four measurement techniques to probe the superconducting gap structure in single-crystal samples of $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$, finding reliable evidence for an isotropic one- or possibly two-gap s-wave model. Thermal conductivity measurements show no sign of low energy quasiparticle excitations even at relatively high magnetic fields, with behavior comparable to other isotropic, single-gapped s-wave superconductors. While there is no indication of a reduced energy gap, the presence of a second gap cannot be excluded based on the small relative field range that was studied.

Conductivity spectra measured by point-contact Andreev reflection spectroscopy exhibit sharp enhancements and notable suppression of dI/dV at lower and higher bias, respectively, which suggests the presence of two gaps. Fitting to an isotropic two-gap BTK model results in gap size estimates of 2.5 meV and 7.0 meV, corresponding to features in the spectra that have been replicated in several crystals from the same batch.

ARPES measurements observe an isotropic gap on both electron and hole bands with a magnitude of approximately 3 meV. Finally, Raman spectroscopy revealed excitations in the superconducting state in both A_{1g} and B_{2g} channels whose energy scales correspond with the 3 meV gap magnitude observed in other measurements.

Overall, we conclude that the optimally doped iron-based superconductor $\text{BaFe}_{1.9}\text{Pt}_{0.1}\text{As}_2$ has an isotropic, fully gapped superconducting order parameter with no nodes or deep minima, possibly with band-dependent energy scales of order 3 meV and 7 meV. Combining several experimental probes together rules out several extrinsic parameters, allowing for further elucidation of the peculiar multi-band nature and pairing mechanism of this and other iron-based superconductors.

Acknowledgments

The authors would like to acknowledge R L Greene and I Takeuchi for valuable discussion and X Zhang for experimental and analytical assistance. Research at the University of Maryland was supported by AFOSR-MURI Grant No. FA9550-09-1-0603, NSF-CAREER Grant No. DMR-0952716, and a travel grant from ICMR. Work at Sherbrooke was supported by the Canadian Institute for Advanced Research and funded by NSERC, CFI, FRQNT, and a Canada Research Chair. Work at Rutgers was supported by the US Department of Energy, Office of Basic Energy Sciences, Division of Materials Sciences and Engineering under Award DE-SC005463.

References

- [1] Kamihara Y *et al* 2008 *J. Am. Chem. Soc.* **130** 3296
- [2] Paglione J and Greene R L 2010 *Nat. Phys.* **6** 645
Johnston D C 2010 *Adv. Phys.* **59** 803
Stewart G R 2011 *Rev. Mod. Phys.* **83** 1589
- [3] Saha S R *et al* 2010 *J. Phys. Condens. Matter* **22** 072204
- [4] Zhu X *et al* 2010 *Phys. Rev. B* **81** 104525
- [5] Hirschfeld P J *et al* 2011 *Rep. Prog. Phys.* **74** 124508
- [6] Taillefer L *et al* 1997 *Phys. Rev. Lett.* **79** 483
- [7] Proust C *et al* 2002 *Phys. Rev. Lett.* **89** 147003
- [8] Reid J-P *et al* 2010 *Phys. Rev. B* **82** 064501
- [9] Reid J-P *et al* 2012 *Phys. Rev. Lett.* **109** 087001
- [10] Tanatar M A *et al* 2011 *Phys. Rev. B* **84** 054507
- [11] Boaknin E *et al* 2003 *Phys. Rev. Lett.* **90** 117003
- [12] Blonder G E *et al* 1982 *Phys. Rev. B* **25** 4515–32
- [13] Daghero D *et al* 2011 *Rep. Prog. Phys.* **74** 124509
- [14] Samuely P *et al* 2009 *Physica C* **469** 507
- [15] Szabó P *et al* 2009 *Phys. Rev. B* **79** 012503

- [16] Ren C *et al* 2011 arXiv:[1106.2891](#)
- [17] Ding H *et al* 2008 *Europhys. Lett.* **83** 47001
- [18] Liu C *et al* 2010 *Nat. Phys.* **6** 419–23
- [19] Terashima K *et al* 2009 *Proc. Natl. Acad. Sci. USA* **106** 7330–3
- [20] Arham H Z and Greene L H 2013 *Curr. Opin. Solid State Mater. Sci.* **17** 81
- [21] Muschler B *et al* 2009 *Phys. Rev. B* **80** 180510
- [22] Saha S R *et al* 2009 *Phys. Rev. Lett.* **103** 037005
- [23] Kirshenbaum K *et al* 2010 *Phys. Rev. B* **82** 144518
- [24] Butch N P *et al* 2010 *Phys. Rev. B* **81** 024518
- [25] Lowell J and Sousa J B 1970 *J. Low Temp. Phys.* **3** 65
- [26] Reid J-Ph *et al* 2012 *Supercond. Sci. Technol.* **25** 084013
- [27] Borisenko S V *et al* 2010 *Phys. Rev. Lett.* **105** 067002
- [28] Umezawa K *et al* 2012 *Phys. Rev. Lett.* **108** 037002
- [29] Allan M P *et al* 2012 *Science* **336** 563
- [30] Andreev A F 1964 *Sov. Phys. JETP* **19** 1228
- [31] Naidyuk Y G *et al* 2002 *JETP Lett.* **75** 238
- [32] Kirshenbaum K *et al* 2012 *Phys. Rev. B* **86** 140505
- [33] Sheet G *et al* 2004 *Phys. Rev. B* **69** 134507
- [34] Lu X *et al* 2010 *Supercond. Sci. Technol.* **23** 054009
- [35] Zhang X *et al* 2010 *Phys. Rev. B* **82** 020515
- [36] Zhang X *et al* 2012 *Phys. Rev. B* **85** 094521
- [37] Hoffman J E 2011 *Rep. Prog. Phys.* **74** 124513
- [38] Yin Y *et al* 2009 *Phys. Rev. Lett.* **102** 097002
- [39] Richard P *et al* 2011 *Rep. Prog. Phys.* **74** 124512
- [40] Nakayama K *et al* 2011 *Phys. Rev. B* **83** 020501
Nakayama K *et al* 2009 *Europhys. Lett.* **85** 67002
- [41] Sauer C and Blumberg G 2010 *Phys. Rev. B* **82** 014525
- [42] Mazin I I *et al* 2010 *Phys. Rev. B* **82** 180502
- [43] Chauvière L *et al* 2010 *Phys. Rev. B* **82** 180521
- [44] Sugai S *et al* 2010 *Phys. Rev. B* **82** 140504